

# Classical and Quantum Nonlinear Dynamics

## Frühjahrssemester 2024

### Blatt 8

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- (1) **Kuramoto's proof that drifters do not contribute** (2 Punkte)  
Kuramoto noted that the stationary distribution of drifting oscillators satisfies the following symmetry:

$$\rho(\theta, \omega) = \rho(\theta + \pi, -\omega)$$

for  $|\omega| > Kr$ . Use that symmetry and the assumption that the density  $g(\omega)$  of natural frequencies is even to prove that

$$\langle e^{i\theta} \rangle_{\text{drift}} = 0$$

in steady state.

- (2) **Square-root growth of the Kuramoto order parameter  $r$  near  $K_c$**  (3 Punkte)  
The goal of this exercise is to show that  $r_\infty(K)$  grows like  $(K - K_c)^{1/2}$  for  $K$  just above  $K_c$  if the unimodal density  $g$  has a quadratic maximum.
- (a) Expand  $g(\omega)$  in a Taylor series for small  $\omega$ . Why is it appropriate to assume that  $\omega$  is small?
- (b) Substitute  $\omega = Kr \sin \theta$  into your Taylor series formula. Explain what this substitution means and why it is valid.
- (c) Assuming that  $r$  is small when  $K$  is near  $K_c$ , show that the self-consistency equation implies

$$r \approx b\sqrt{K - K_c}$$

and find the prefactor  $b$  explicitly in terms of  $K_c$  and  $g''(0)$ .

- (3) **Locking threshold for parabolic density of natural frequencies** (5 Punkte)  
Consider the following density for the natural frequencies in the Kuramoto model:

$$g(\omega) = \frac{3}{4}(1 - \omega^2)$$

for  $-1 \leq \omega \leq 1$ , and  $g(\omega) = 0$  otherwise. In cases like this where the density does not have infinite tails, the Kuramoto model undergoes *two* transitions: a transition from incoherence to partial locking at  $K = K_c$  followed by a transition from partial locking to complete locking at  $K = K_L \geq K_c$ .

- (a) Determine  $K_c$  and  $K_L$  by choosing appropriate values of  $Kr$  in the integrand of the self-consistency equation.

Hint for  $K_L$ : what is the minimal value of  $Kr$  such that all phases will lock?

- (b) Show that the order parameter for the partially-locked states is given by

$$r_\infty(K) = \frac{2}{K} \sqrt{1 - \frac{8}{3\pi K}}.$$

- (c) For  $K > K_L$  the system is completely locked. The phases of the locked oscillators no longer span the whole semicircle  $|\theta| \leq \pi/2$ , but squeeze together more tightly. Find the maximum locked phase  $\theta_{\max}$ .

- (d) (Bonus problem) To compute  $r_\infty(K)$  for the completely locked states, we can use the parametric equation method of problem 13.5.5 in Strogatz.

Define  $u = Kr$ , rewrite  $\theta_{\max}$  in terms of  $u$  and modify the limits of integration in the self-consistency equation accordingly. Then do the same for  $f(u)$ . Show that

$$f(u) = \frac{3}{16} \left( \frac{\sqrt{u^2 - 1}(u^2 + 2)}{u^2} - (u^2 - 4) \csc^{-1}(u) \right).$$

Find  $r(u)$  and  $K(u)$  and use them to plot  $r(K)$  for the completely-locked states.

- (e) (Bonus problem) As an alternative to (d), simulate the Kuramoto model numerically for the parabolic density and try to estimate  $K_c$  and  $K_L$ .